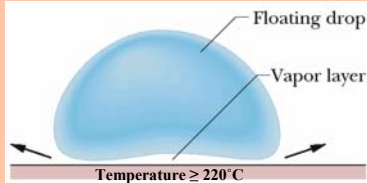


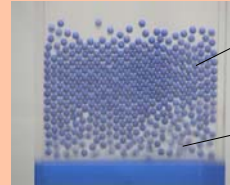
Peter Eshuis, Ko van der Weele, Devaraj van der Meer, Detlef Lohse

## Johann Gottlob Leidenfrost (1756)



A drop of water hovering over a hot plate on its own vapor layer.

## Granular Leidenfrost effect



Floating cluster

Vapor-like layer



Shaker and flat container with blue beads

Vertically vibrated glass beads: Above a critical shaking strength, a crystalline cluster (analogous to the Leidenfrost drop) floats on a vapor-like layer of much faster beads.

## How to distinguish the cluster from the vapor-like region?

We employ the concept of pair correlations between the particles ( $N$ ) in a horizontal strip ( $y, y+dy$ ):

$$g_y(x) = \frac{1}{N} \sum_{i,j \text{ in } (y,y+dy)} \sum_{i \neq j} \delta(x - (x_i - x_j))$$

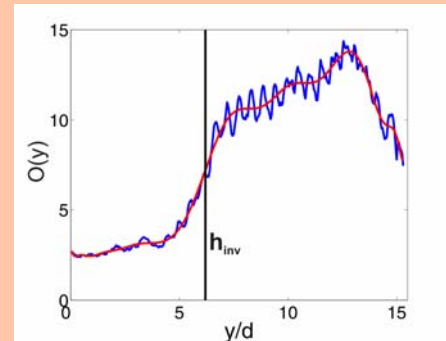
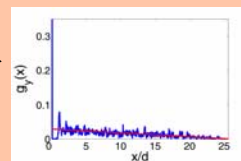
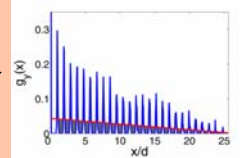
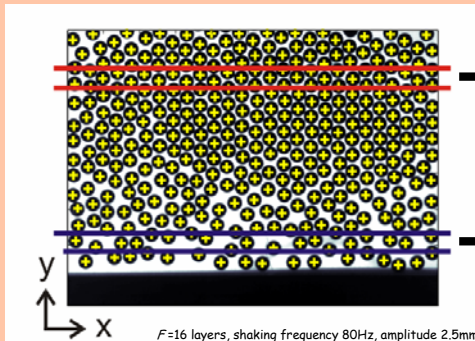
- Strip near bottom: particle positions show hardly any correlation, i.e., typically gas-like.
- Strip inside floating cluster: strong periodic, crystalline correlation in the  $x$  direction.

This clear distinction between periodic and non-periodic behavior in  $g_y(x)$  is used in the order parameter  $O$ :

$$O(y) = \int |g_y(x) - g_{y,fit}(x)| dx \quad (\text{the shaded area in the pair correlation plots})$$

The transition from gas to crystal (the so-called inversion height  $h_{inv}$ ) is located at the point where the slope of the fit through  $O(y)$  is maximal.

A non-zero inversion height is our key indicator for the Leidenfrost effect.



## Hydrodynamic model

The effect is quantitatively described by three hydrodynamic equations:

Equation of state  $p = nT \frac{n_c + n}{n_c - n}$

Force balance  $\frac{dp}{dy} = -mgn$

Energy balance  $\frac{d}{dy} \left( \kappa \frac{dT}{dy} \right) = \frac{\mu}{\gamma l} \varepsilon n T^{\frac{3}{2}}$   
(heat flux) (dissipation)

$p$  - pressure  
 $n$  - number density ( $n_c = 2/\sqrt{3}d^2$ )  
 $T$  - granular temperature  
 $m$  - mass of a single particle  
 $\kappa$  - thermal conductivity  
 $l$  - mean free path  
 $\mu, \gamma$  - constants

Non-dimensionalizing these equations (and the associated boundary conditions) we find that the problem is governed by **three dimensionless control parameters**:

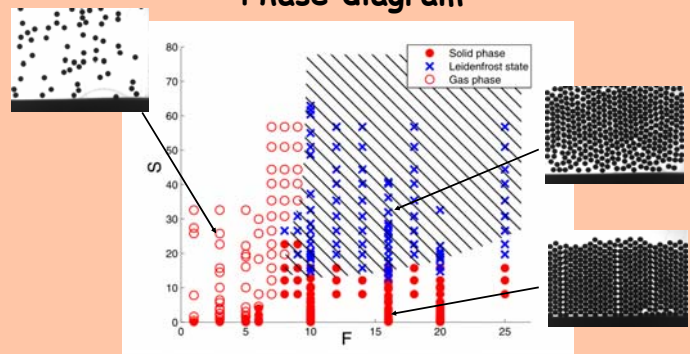
[1] The energy input  $S = 4\pi^2 \frac{m(ad)^2}{mgd} = \frac{\text{typical kinetic energy a bead gets from bottom}}{\text{energy needed for a bead to jump to height } d}$

[2] The inelasticity  $\varepsilon = (1 - e^2)$

[3] Number of particle layers  $F$

$d$  - particle diameter  
 $e$  - coefficient of restitution

## Phase diagram



- Excellent agreement between experiment (dots, crosses) and theory (shaded area).
- The Leidenfrost state shows up between the solid and gas phase, for sufficiently large shaking strength ( $S \geq 16$ ) and number of layers ( $F \geq 8$ ).

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